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Electronic and Magnetic Properties of π -d Interaction System: (EDTDM)₂FeBr₄

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(EDTDM)₂FeBr₄ based on π -d interaction system consists of quasi one-dimensional donor columns and square-lattice magnetic anion sheets (Fe³⁺(s=5/2)). It behaves metallic at high temperatures and undergoes a metal-insulator transition at $T_{\rm MI}\sim15$ -20K, where the low temperature insulating phase is considered to be magnetic ordered state. The magnetic susceptibility, which is governed by Fe³⁺ d-spins, obeys the Curie-Weiss law in the high-temperature regime. Below 50K, the susceptibility deviates from the Curie-Weiss law and shows an antiferromagnetic transition at $T_{\rm N}\sim3$ K. The anomalous magnetization curves indicate an important role of π -d interaction in the magnetic behavior.

<u>Keywords</u>: TTF derivative; M-I transition; antiferromagnetic transition; π -d interaction

INTRODUCTION

Organic donor molecules, such as TTF derivatives, form a large variety of charge transfer complexes with acceptors, in which planar-shaped donor molecules form one- or two-dimesional arrangements through a face-to-face and/or side-by-side inter-molecular interactions. Consequently, the overlap of the π -electron wave functions between adjacent donors produce a low dimensional electronic structures, which are stabilized in a metallic or semiconducting phase depending on charge transfer rate, manner of molecular arrangement, Coulomb interactions, etc. When we employ transition metal anions that have localized d-electron spins as a constituent, the d-electron

spins are expected to interact with donor π -electrons in the complexes. This π -d interaction will produce metal magnetism in the case of metallic π -electron system or π -d composite magnetism in the semiconducting phase [1, 2, 3].

Recently, Enomoto et al. [4] reported that $(DMET)_2FeBr_4$ shows the interesting interplay between magnetism and electron transport, which is associated with the interaction between the localized delectrons and the conduction π -electrons. If the donor molecules of this system can be substituted for analogues with keeping the crystal structure unchanged, it is possible to investigate systematically the role of donor molecules in the π -d interaction system.

In this study, we focused on EDTDM (FIGURE 1) as a donor molecule where sulfur is substituted for selenium in DMET molecule, and the crystal structure and physical properties of (EDTDM)₂FeBr₄ and (EDTDM)₂GaBr₄ are investigated.

FIGURE 1: The molecular structures of EDTDM (left) and DMET (right) .

EXPERIMENTAL

Single crystals of (EDTDM)₂FeBr₄ were obtained by a galvanostatic anodic oxidation method (0.1-0.5 μ A) for a synthetic period of 2 weeks using 9mg of EDTDM, 25mg of TBAFeBr₄ as a supporting electrolyte and 15cm³ of chlorobenzene as a solvent. Single crystals of (EDTDM)₂GaBr₄ were obtained by using TBAGaBr₄ under the same condition. The crystal structures were determined by a single crystal X-ray diffraction method using RIGAKU AFC-7 four-circle diffractometer. The reflection data were collected in the range of 2θ =5 to 55°. Their structures were solved with the direct methods using SHELXS-86, and then refined with the full-matrix least-squares method using SHELXL-93. The coordinates of hydrogen atoms were calculated geometrically. The magnetic susceptibilities

were measured using Quantum-Design MPMS-5 SQUID magnetometer up to a field of 5.5T in the temperature range 1.85-300K. Electrical resistivity was measured by constant current four-probe method ($R \le 10k\Omega$), or constant voltage two-probe method ($R \ge 10k\Omega$). ESR measurements were performed with an ESR X-band spectrometer (JEOL JES-TE200) equipped with a continuous-flow cryostat (Oxford ESR910) in the temperature range 3.3K-300K.

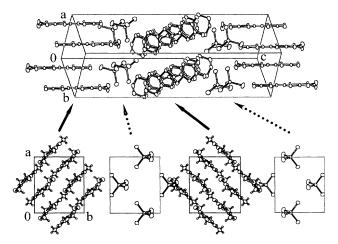
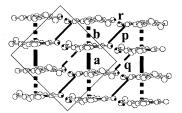


FIGURE 2: The crystal structure of (EDTDM)₂FeBr₄.

RESULTS AND DISCUSSION

Both of the obtained salts, $(EDTDM)_2FeBr_4$ and $(EDTDM)_2GaBr_4$, are isostructural to $(DMET)_2FeBr_4$ [4] with a 2:1 donor/anion composition. The crystal structure of $(EDTDM)_2FeBr_4$ is shown in FIGURE 2, which is characterized with an alternating stacking of EDTDM conducting sheets and $FeBr_4^-$ magnetic sheets. EDTDM columns are arranged in a face-to-face manner in the ab-plane where the columns on the adjacent layers are directed two different directions; one is along the a+b direction and the other is along the a-b direction. Two columns are crystallographycally equivalent, for the asymetric unit contains one molecule for donors and a half for anions.

From comparision of the intra-chain overlap integrals between (EDTDM)₂FeBr₄ and (DMET)₂FeBr₄ shown in FIGURE 3, it is re-



site	EDTDM	DMET
$\mathbf{a}(\mathbf{a}/\mathbf{b})$	24.92(1.12)	60.52(1.02)
$\mathbf{b}(\mathbf{b}/\mathbf{b})$	22.19(1)	59.55(1)
$\mathbf{p}(\mathbf{p}/\mathbf{b})$	1.502(0.0677)	10.81(0.182)
$\mathbf{q}(\mathbf{q}/\mathbf{b})$	2.787(0.126)	6.742(0.113)
$\mathbf{r}(\mathbf{r}/\mathbf{b})$	2.360(0.106)	6.237 (0.105)

FIGURE 3: The intra-chain (\mathbf{a}, \mathbf{b}) and inter-chain $(\mathbf{p}, \mathbf{q}, \mathbf{r})$ overlap integrals of $(EDTDM)_2FeBr_4$ and $(DMET)_2FeBr_4$ [4] calculated from extended Hückel approximation. The values of the overlap integrals that are normalized by the value of \mathbf{b} are shown in parenthese.

vealed that the dimerization of donor molecules in (EDTDM)₂FeBr₄ is a little stronger than that in (DMET)₂FeBr₄. In addition, it is indicated that donor arrangement for (EDTDM)₂FeBr₄ is more one-dimensional than that for (DMET)₂FeBr₄. Actually, the intra-chain transfer integrals are five times larger than the inter-chain transfer integrals in (EDTDM)₂FeBr₄.

The anion molecules form two-dimensional square lattice in the ab-plane. The inter-anion Br-Br distance is estimated at 3.87Å for (EDTDM)₂FeBr₄ and 3.88Å for (EDTDM)₂GaBr₄, which is comparable with the sum of van der Waals radii (3.7Å), suggesting the presence of the inter-anion exchange interaction paths of the close inter-anion Br-Br distance. Moreover, there are close contacts between an anion and its neighboring donor molecules, whose intermolecular S-Br distance is estimated at 3.74Å for (EDTDM)₂FeBr₄ and 3.76Å for (EDTDM)₂GaBr₄, comparable to the sum of van der Waals radii (3.65Å). Therefore, we cannot ignore the interaction between adjacent two anion sheets, which are bridged by the donor columns

For both (EDTDM)₂FeBr₄ and (EDTDM)₂GaBr₄ the in-plane resistivity shows metallic behavior above $\sim 200 \mathrm{K}$ and apparently semiconducting behavior below $\sim 200 \mathrm{K}$. In contrast, the out-of-plane resistivity shows semiconducting behavior for both crystals. Therefore, both (EDTDM)₂FeBr₄ and (EDTDM)₂GaBr₄ have the similar transport property.

The magnetic susceptibility of $(EDTDM)_2GaBr_4$ shows Pauli paramagnetic behavior above $\sim 20K$, whose contribution gradually de-

creases as the temperature decreases. Since the Pauli paramagnetic susceptibility is proportional to the conduction carrier density, the decrease in the density of states explains the apparent semiconductor-like behavior in the resistivity for $T \le 200$ K even in the metallic phase though the presence of Fermi surface is predicted from band structure calculation at room temperature. Below ~ 20 K the susceptibility decreases steeply as the temperature decreases, suggesting the appearance of an antiferromagnetic (AF) ordered state, taking into account NMR measurement [5].

The ESR signal of (EDTDM)₂GaBr₄ appears as one Lorentzian peak with line-width $\Delta H_{\rm pp} \sim 0.7 \,\mathrm{mT}$ and g-value g ~ 2.01 at room temperature when the field is applied along the c^* -axis. In contrast, in the applied field along the a+b or the a-b direction, two Lorentzian peaks are observed with line-width $\Delta H_{\rm pp} \sim 0.65 \,\mathrm{mT}$ and $0.75 \,\mathrm{mT}$ and g-value $g\sim 2.0079$ and 2.0025. The presence of two peaks is the consequence of 'double-column' structure. Since the long axes of donor molecules are nearly parallel to the c-axis, two donor columns are equivalent for $B//c^*$ -axis. On the other hand, when the field is applied in the conducting plane, two columns are generally inequivalent, and giving double peaks in the ESR spectra. The ESR line-width $\Delta H_{\rm pp}$ under the field along c*-axis decreases with decreasing temperature and below ~ 15 K the line-width abruptly diverges. The intensity has slight temperature dependence above ~ 15 K and rapidly decreases below ~ 15 K. The temperature dependence is consistent with the appearance of AF state below 15K, as the susceptibility and NMR suggest. SDW or Mott-like AF state is a plausible candidate of the ordered state below $T_{\rm MI}$ =15-20K.

FIGURE 4 show the temperature dependence of the magnetic susceptibility in an applied field of 1T (a) and the magnetization curves (b) for (EDTDM)₂FeBr₄. The susceptibility obeys the Curie-Weiss law in the high temperature region above ~ 50 K, where the Curie constant C and the Weiss temperature Θ are estimated at C=4.74emu·K/mol and $\Theta=-2.8$ K (200K<T<300K), respectively. From the Curie constant, it is proved that Fe³⁺ spin (S=5/2) mainly serves as the origin of the magnetism of (EDTDM)₂FeBr₄. This salt shows antiferromagnetic transition at $T_{\rm N}\sim 3$ K, below which the magnetic easy-axis is the a-axis. In addition, anisotropy appears in the susceptibility below ~ 50 K, and it increases with decreasing the temperature. From the magnetization curves, a spin-flop transition

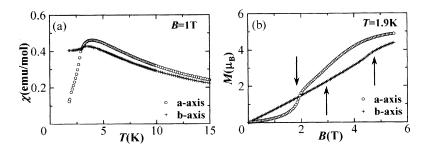


FIGURE 4: The temperature dependence of the magnetic susceptibility (a) and the magnetization curves (b) of (EDTDM)₂FeBr₄.

occurs at $B_{\rm sf} \sim 1.8 {\rm T}$ in addition to the appearance of an anomaly at $\sim 2.9 {\rm T}$ when the magnetic field is applied along the a-axis. For B//b-axis, the magnetization curve has a shoulder around $B \sim 4.5 {\rm T}$. A large difference emerges in the magnetizations between B//a-axis and B//b-axis. These unusual behaviors of the magnetization cannot be explained only by Fe³⁺ spins. The donor spins is considered to participate in the magnetization as in the case of (DMET)₂FeBr₄ [4].

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